

## Research Article

# Investigation on Domain Pinning Mechanism in Nanometer Hard Magnetic Materials

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Received 5 March 2019; Revised 11 June 2019; Accepted 24 June 2019; Published 6 August 2019

Academic Editor: Oleg Derzhko

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The domain pinning mechanism was investigated in nanometer hard magnetic materials. The pinning fields of domain wall at different inhomogeneities were studied respectively. And exchange-coupling coefficient  $\alpha_{ex}$  was investigated too. The results showed that  $\alpha_{ex}$  is proportional to the ratio of inhomogeneity thickness  $r_o$  to wall width  $\delta_m$  for narrow inhomogeneity, while it decreases with enhancement of  $r_o/\delta_m$  for extended inhomogeneity. At a certain value of  $r_o$ , exchange-coupling coefficient and pinning fields will reach the maximum. Exchange-coupling interaction and pinning fields are greatly influenced by inhomogeneity. Control range of inhomogeneity may obtain higher coercivity.

## 1. Introduction

In recent years, nanocomposite magnets have attracted much attention due to high theoretical energy product, which was recorded as high as  $1 \text{ MJ/m}^3$  [1]. However, an equally important uncertainty is coercivity mechanism in nanometer magnetic materials. Therefore, attention has been focused on coercivity mechanism [2–5]. Currently, there are mainly two coercivity mechanisms, nucleation and pinning. Kronmüller et al. [5] proposed that the coercivity of nanocomposite magnets was controlled by nucleation process of reverse domains and introduced a microstructural parameter  $\alpha_{ex}$  to describe the effect of exchange-coupling interaction on coercivity, leading to  $\mu_0 H_c = \alpha_k \alpha_{ex} \mu_0 H_N^{\min} - N_{eff} J_s$  (coercivity experiential formula). Zhang et al. [6] suggested both nucleation and domain-wall pinning model can be expressed as coercivity experiential formula. Experimentally, elemental additions were found to form precipitates at grain boundary [7, 8]. The properties of grain boundary are different from main phase. Therefore, there is inhomogeneous area in nanometer magnetic materials, which plays an important role in pinning domain wall. In this paper, pinning fields of domain wall  $H_c$  were calculated by considering inhomogeneous area. And corresponding exchange-coupling coefficient  $\alpha_{ex}$  was investigated too.

## 2. Calculation Model

**2.1. Inhomogeneous Area.** Experimentally, elemental additions were found to form precipitates at grain boundary [7, 8]. Properties of grain boundary are different from the main phase. Additionally, due to the exchange-coupling interaction between grains, properties of grain edge have changed too. It is assumed that grain boundary and area influenced by exchange-coupling interaction form inhomogeneous area. Referring to  $K_1$ -profile given by Kronmüller [9], we assume anisotropy at inhomogeneous area is as follows:

$$K_1(r) = K_1 - \frac{nK_1}{ch^2(r/r_o)}, \quad (1)$$

where  $r_o$  is inhomogeneity thickness.  $K_1$  is first anisotropy constant.  $n$  is a number no larger than 1, which denotes reduction  $K_1$  in inhomogeneous area. The variation of anisotropy at inhomogeneous area is shown in Figure 1.

**2.2. Coercivity Theory.** Inhomogeneous area is regarded as strong planar pinning centers. When inhomogeneity thickness  $r_o$  is smaller than wall width  $\delta_m$ , coercive field is given [10]:

$$H_c = \frac{1}{3\sqrt{3}} \frac{2K_1}{M_s} \frac{\pi d}{\delta_m} \sum_1^n \left( \frac{A}{A^{i,i+1}} - \frac{K_1^i}{K_1} \right) - N_{eff} M_s, \quad (2)$$

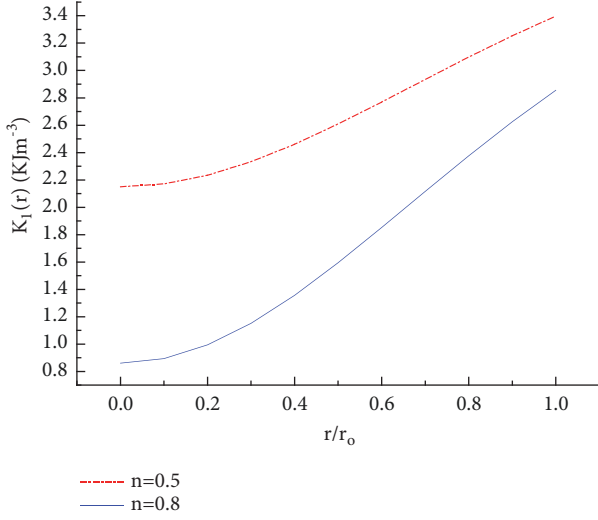


FIGURE 1: Variations of anisotropy  $K_1(r)$  in inhomogeneous area.

where  $A^{i,i+1}$  corresponds to the local exchange constant between adjacent layers  $i$  and  $i+1$ . Similarly,  $K_1^i$  corresponds to the local anisotropy constant of the  $i$ th layer.  $\delta_m$  is the wall width. The term  $N_{eff}M_s$  takes care of demagnetization fields resulting from grain surfaces and volume charges. It is assumed that the  $n$  perturbed planes have equivalent properties ( $A'$  and  $K_1'$  are equal). And we assume  $A'$  of perturbed planes is equal to  $A$ . Coercivity is written as

$$H_c = \frac{1}{3\sqrt{3}} \frac{2K_1}{M_s} \frac{\pi r_o}{\delta m} \left( 1 - \frac{K_1'}{K_1} \right) - N_{eff}M_s, \quad (3)$$

where  $r_o = md$  denoting inhomogeneous area thickness.

When  $r_o$  is larger than  $\delta_m$ , the coercive field is determined by

$$H_c^{ij} = \frac{1}{2M_s} \left( \frac{d\gamma}{dr} \right)_{\max} - N_{eff}M_s, \quad (4)$$

where  $(d\gamma/dr)_{\max}$  denotes maximum slope of the wall energy given by  $\gamma = 4\sqrt{A_1K_1(r)}$ .  $K_1(r)$  is anisotropy at inhomogeneous area.

**2.3. Coefficient for Exchange-Coupling Interaction.** Compared with coercivity experiential formula, expressions of exchange-coupling coefficient  $\alpha_{ex}$  are given. For thin inhomogeneity ( $r_o < \delta_m$ ), exchange-coupling coefficient  $\alpha_{ex}$  is written as

$$\alpha_{ex} = \frac{1}{3\sqrt{3}} \frac{\pi r_o}{\delta_B} \left( 1 - \frac{K_1'}{K_1} \right). \quad (5)$$

In this case, for  $K_1' = K_1(r_o/2)$  and  $n=0.8$ ,  $\alpha_{ex} = (7\sqrt{3}/100)(\pi r_o/\delta_B)$  is obtained.

For thick inhomogeneity ( $r_o > \delta_m$ ),  $\alpha_{ex}$  is written as

$$\alpha_{ex} = \frac{n\delta_m}{\pi r_o} \left( \frac{d\sqrt{K_1(r)}}{dr} \right)_{\max}. \quad (6)$$

And when  $n$  equals 0.8,  $\alpha_{ex} = 12\delta_m/25\pi r_o$ .

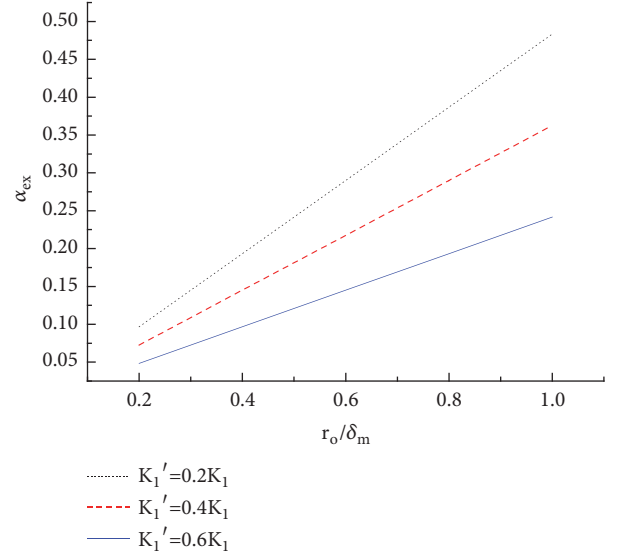


FIGURE 2: Variations of exchange-coupling coefficient  $\alpha_{ex}$  with the ratio of  $r_o$  to  $\delta_m$  for narrow inhomogeneity.

### 3. Results and Discussion

We calculated coercivity and corresponding  $\alpha_{ex}$  by using intrinsic magnetic parameters of  $\text{Nd}_2\text{Fe}_{14}\text{B}$  ( $K_1=4.3\text{MJ/m}^3$ [5],  $J=1.6\text{T}$ ,  $M_s=1280\text{KA/m}$ [11]).

Figure 1 shows variations of anisotropy  $K_1(r)$  in inhomogeneous area. It can be seen that  $K_1(r)$  increases with enhancement of ratio of  $r$  to  $r_o$ . The value of anisotropy of area close to grain core is larger, because it is influenced to a fewer degree by exchange-coupling interaction. It is also shown in Figure 1 that  $K_1(r)$  decreases with increasing  $n$ .  $n$  denotes the decrease  $K_1$  in inhomogeneous area. Exchange-coupling interaction and inhomogeneity decrease anisotropy.

Figure 2 shows variations of exchange-coupling coefficient  $\alpha_{ex}$  with ratio of  $r_o$  to  $\delta_m$  for narrow inhomogeneity. It can be seen that  $\alpha_{ex}$  is proportional to  $r_o/\delta_m$  in this case. For smaller  $r_o$ , exchange-coupling interaction is strong, and the effects of exchange-coupling interaction on anisotropy and coercivity are dominant. The enhancement of  $\alpha_{ex}$  mainly results from influence of exchange-coupling interaction. Figure 3 shows variations of exchange-coupling coefficient  $\alpha_{ex}$  with ratio of  $r_o$  to  $\delta_m$  for extended inhomogeneity. It can be seen that  $\alpha_{ex}$  decreases with increasing of  $r_o/\delta_m$  for certain value of  $n$ . For extended inhomogeneity, the effect of defect on anisotropy and coercivity is dominant.

Figure 4 shows dependence of  $\alpha_{ex}$  on ratio of  $r_o$  to  $\delta_m$ . In transition region near 0.6, the broken curve indicates the transition between the approximations for narrow and extended inhomogeneities. As shown in Figure 5, the variation of coercivity and exchange-coupling coefficient is consistent. Inhomogeneity has great influence on coercivity. Magnetization reversal process is thought of as follows. Firstly, the antimagnetization core is formed in the nonuniform region. With enhancement of the antimagnetization field, domain wall displacement will encounter maximum resistance peak caused by exchange coupling interaction

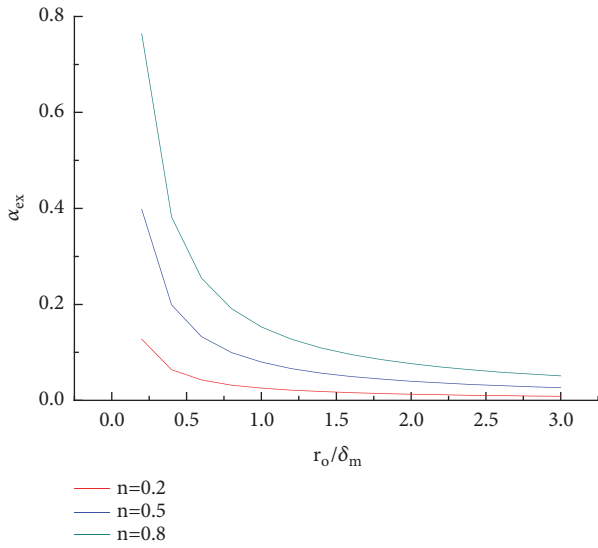


FIGURE 3: Variation of exchange-coupling coefficient  $\alpha_{ex}$  with the ratio of  $r_o$  to  $\delta_m$  for extended inhomogeneity.

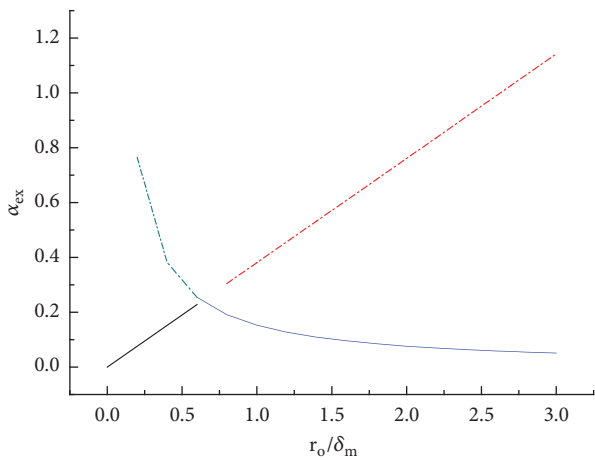


FIGURE 4: Dependence of  $\alpha_{ex}$  on ratio of  $r_o$  to  $\delta_m$ .

and inhomogeneity, and domain wall will have maximum irreversible displacement. Pinning field reaches the maximum value. With increasing inhomogeneity, area influenced by exchange-coupling interaction is larger for narrow inhomogeneity, while area influenced by exchange-coupling interaction is smaller with increasing range of inhomogeneity for extended inhomogeneity.

#### 4. Conclusion

It is assumed that defect and grain edge influenced by exchange-coupling interaction form inhomogeneous area. Pinning fields and exchange-coupling coefficient  $\alpha_{ex}$  corresponding to narrow and extended inhomogeneous area are investigated respectively. By considering different inhomogeneities, pinning fields and exchange-coupling coefficient were studied. While the effect of defect on them is dominant

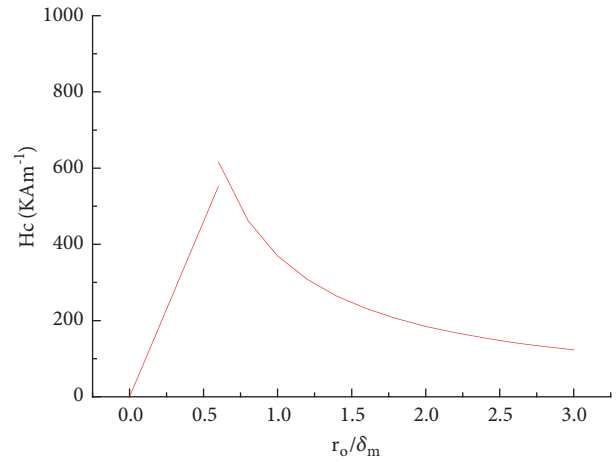


FIGURE 5: Dependence of coercivity on ratio of  $r_o$  to  $\delta_m$ .

for extended inhomogeneity. Coercivity depends greatly on microstructure.

#### Data Availability

The data used to support the findings of this study are included within the article.

#### Conflicts of Interest

The authors declare that they have no conflicts of interest.

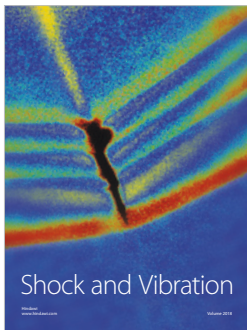
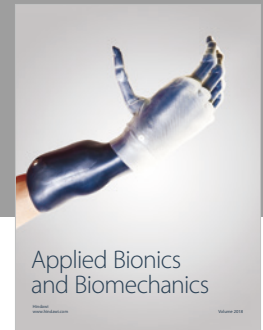
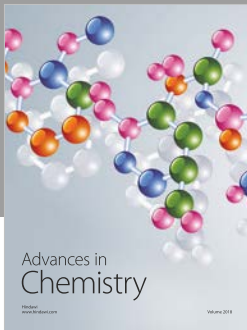
#### Acknowledgments

The work is supported by the Natural Science Foundation of Shandong Province (ZR2014EL002).

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